

MPHYCC10-13

Quantum mechanical relativity correction:

Relativistic effect is equally important as it produces energy shift comparable to spin-orbit interaction (in Hydrogen atom)

$$H = K + V \quad \dots \quad (1)$$

$K$  = relativistic kinetic energy,  $V$  = potential energy

$$K = (p^2 c^2 + m_0^2 c^4)^{1/2} - m_0 c^2 \quad p = \text{momentum } m_0 = \text{rest mass of electron.}$$

$H$  = Hamiltonian

$$H = (p^2 c^2 + m_0^2 c^4)^{1/2} - m_0 c^2 + V$$

$$= m_0 c^2 \left[ 1 + \frac{p^2}{m_0^2 c^2} \right]^{1/2} - m_0 c^2 + V$$

$$= m_0 c^2 \left[ 1 + \frac{1}{2} \frac{p^2}{m_0^2 c^2} - \frac{1}{8} \frac{p^4}{m_0^4 c^4} + \dots \right] - m_0 c^2 + V$$

$$= m_0 c^2 + \frac{p^2}{2 m_0} - \frac{p^4}{8 m_0^3 c^2} + \dots - m_0 c^2 + V$$

$$H = \frac{p^2}{2 m_0} - \frac{p^4}{8 m_0^3 c^2} + \dots + V \quad \dots \quad (2)$$

change in Hamiltonian due to relativity =  $-\frac{p^4}{8 m_0^3 c^2}$  = perturbation term.

$$p = -i \hbar \frac{\partial}{\partial r} = -\frac{i \hbar}{2\pi} \frac{\partial}{\partial r}$$

Then the Hamiltonian operator can be written as

$$H = \frac{1}{8 m_0^3 c^2} \left( -\frac{i \hbar}{2\pi} \frac{\partial}{\partial r} \right)^4 = \frac{1}{8 m_0^3 c^2} \frac{\hbar^4}{16 \pi^4} \nabla^4$$

MPHYCC10-13

If  $\psi_0$  = unperturbed wave function of the hydrogen atom then first order energy shift due to relativity

$$\Delta E_n = - \int \psi_0^* \left[ \frac{h^4 / 16 \pi^4}{8 m_0^3 c^2} \right] \psi_0 d\tau$$

$$\Delta E_n = - \frac{R_{\infty} \alpha^2 Z^4 h c}{n^3} \left[ \frac{1}{l + \frac{1}{2}} - \frac{3}{4n} \right]$$

The relativistic term shift  $\Delta T_n = - \frac{\Delta E_n}{h c} = \frac{R_{\infty} \alpha^2 Z^4}{n^3} \left[ \frac{1}{l + \frac{1}{2}} - \frac{3}{4n} \right]$

Hydrogen fine structure →

First order perturbation corrections due to different effect combine linearly. Hence net term shift due to spin orbit interaction and the relativistic effect in a hydrogen like atom is

$$\Delta T = \Delta T_{l,s} + \Delta T_n$$

$$= - \frac{R_{\infty} \alpha^2 Z^4}{2 n^3 l(l + \frac{1}{2})(l + 1)} [j(j+1) - l(l+1) - s(s+1)] + \frac{R_{\infty} \alpha^2 Z^4}{n^3} \left[ \frac{1}{l + \frac{1}{2}} - \frac{3}{4n} \right]$$

$$= \frac{R_{\infty} \alpha^2 Z^4}{n^3} \left[ \frac{1}{l + \frac{1}{2}} - \frac{j(j+1) - l(l+1) - s(s+1)}{2 l(l + \frac{1}{2})(l + 1)} - \frac{3}{4n} \right] \dots \textcircled{1}$$

$s = \frac{1}{2}$ ,  $J = l + \frac{1}{2}$  &  $J = l - \frac{1}{2}$  we will have

$$j(j+1) - l(l+1) - s(s+1) = l \quad \text{for } J = l + \frac{1}{2} \text{ \& } s = \frac{1}{2} \textcircled{2}$$

$$= -(l+1) \quad \text{for } J = l - \frac{1}{2} \text{ \& } s = \frac{1}{2}$$

using  $\textcircled{2}$  in  $\textcircled{1}$  the net term shifts corresponding to  $j = l + \frac{1}{2}$  &

$J = l - \frac{1}{2}$

$$\Delta T = \frac{R_{\infty} \alpha^2 Z^4}{n^3} \left[ \frac{1}{l + 1} - \frac{3}{4n} \right] \dots \textcircled{3}$$



Dr. Ranjana Singh  
Assistant Professor Physics  
Harshnada Das Jain College  
Bhoga, Bihar, India

(03)

Date: 19/02/2025  
Topic: O.P.M.

MPHYEC10-13

$$\Delta T' = \frac{R_{\infty} \alpha^2 z^4}{n^3} \left[ \frac{1}{l+1} - \frac{3}{4n} \right] \quad \text{--- (4)}$$

$$\Delta T'' = \frac{R_{\infty} \alpha^2 z^4}{n^3} \left[ \frac{1}{l} - \frac{3}{4n} \right] \quad \text{--- (5)} \quad J = l + \frac{1}{2} \quad 4J = 4l + 2$$

then eqn (4) & (5) can be represented by single eqn

$$\Delta T = \frac{R_{\infty} \alpha^2 z^4}{n^3} \left[ \frac{1}{J + \frac{1}{2}} - \frac{3}{4n} \right]$$

$R_{\infty}$  = Rydberg's constant for heavy nucleus =  $1.097 \times 10^7 \text{ m}^{-1}$   $\alpha = \frac{1}{137}$   
 $z=1$  for hydrogen atom.

$$\Delta T = \frac{584}{n^3} \left[ \frac{1}{k} - \frac{3}{4n} \right] \text{cm}^{-1} = \frac{5.84}{n^3} \left[ \frac{1}{k} - \frac{3}{4n} \right] \text{cm}^{-1} \quad \text{Sommerfeld Formula}$$

$$\Delta T = \frac{584}{n^3} \left[ \frac{1}{J + \frac{1}{2}} - \frac{3}{4n} \right] \text{cm}^{-1} = \frac{5.84}{n^3} \left[ \frac{1}{J + \frac{1}{2}} - \frac{3}{4n} \right] \text{cm}^{-1} \quad \text{Dirac Formula}$$

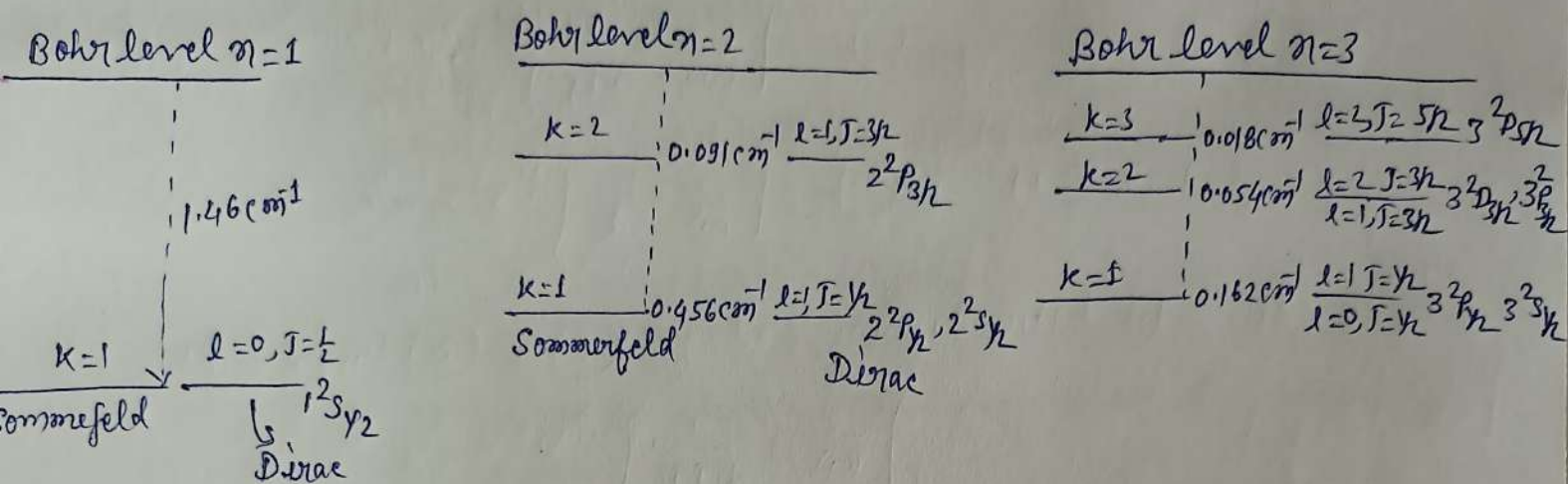
Bohr level $n$	Sommerfeld levels		Dirac levels		$\Delta T (\text{cm}^{-1})$
	$k$	$\Delta T (\text{cm}^{-1})$	$l$	$J = l \pm \frac{1}{2}$	
1	1	1.46	0	$\frac{1}{2}$	1.46
2	2	0.091	1	$\frac{3}{2}, \frac{1}{2}$	0.091, 0.456
	1	0.456	0	<del><math>\frac{3}{2}</math></del> , $\frac{1}{2}$	0.456
3	3	0.018	2	$\frac{5}{2}, \frac{3}{2}$	0.018, 0.054
	2	0.054	1	$\frac{3}{2}, \frac{1}{2}$	0.054, 0.162
	1	0.162	0	$\frac{1}{2}$	0.162

Dr. Ranjana Singh  
 Assistant Professor, Physics  
 Man Prasad Das Jain College  
 Ara, Bihar, India

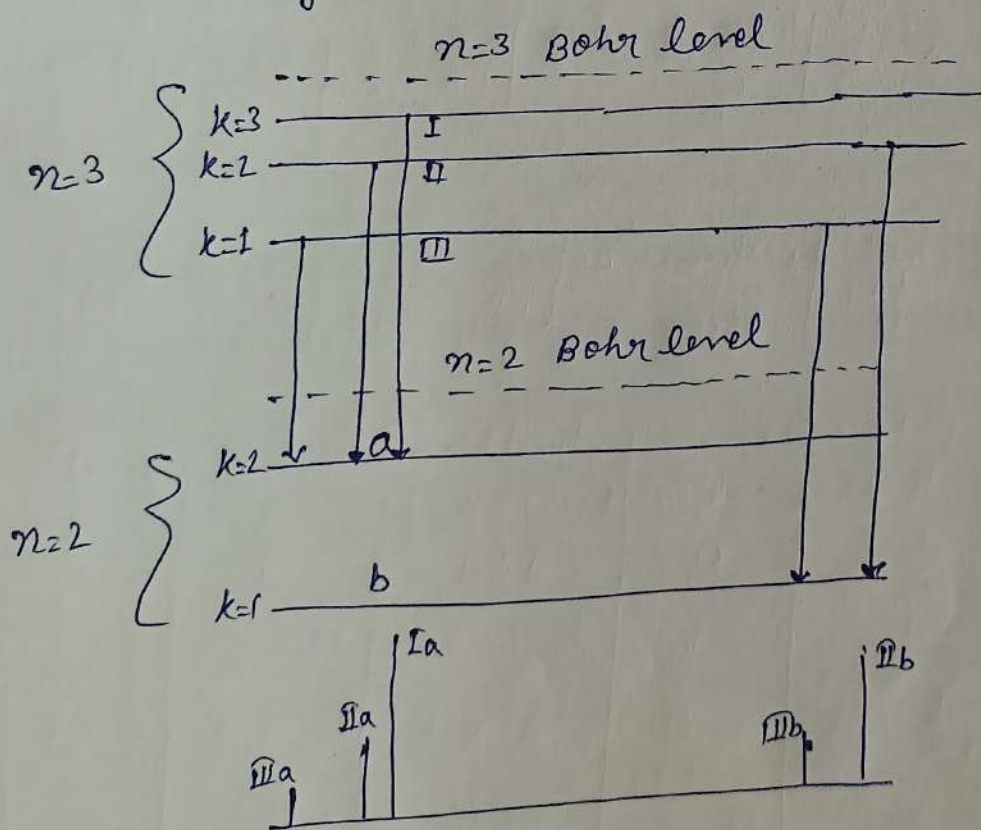
(64)

Date: 19/02/2025  
 Time: 01 pm

MPHYCC10-13



According to Dirac theory most levels are really double  
 Fine structure of  $H\alpha$  line ( $3 \rightarrow 2$ ) on the basis of Dirac theory



The selection rule for the dipole transitions are  
 $\Delta l = \pm 1$   
 $\Delta J = 0, \pm 1$   
 but  $J=0 \not\leftrightarrow J=0$

Theoretically deduced structure of  $H\alpha$  line has general agreement with Henson's observations. still it is away from perfect agreement